Correlated Brownian motion and diffusion of defects in spatially extended chaotic systems

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Resumo

One of the spatiotemporal patterns exhibited by coupled map lattices with nearest-neighbor coupling is the appearance of chaotic defects, which are spatially localized regions of chaotic motion with a particle-like behavior. Chaotic defects display random behavior and diffuse along the lattice with a Gaussian behavior. In this note we investigate some dynamical properties of chaotic defects in a lattice of coupled chaotic quadratic maps. Using a recurrence-based diagnostic we found that the motion of chaotic defects is well-represented by a stochastic time series with a power-law spectrum, i.e. a correlated Brownian motion. The correlation exponent corresponds to a memory effect in the Brownian motion, and increases with a system parameter as the diffusion coefficient of chaotic defects.

Mathematical Section

We use as a paradigmatic model of spatio-temporal dynamics the following coupled map lattice

(i)
$$\epsilon(i) \in \epsilon(i)$$
 $\epsilon(i) \in \epsilon(i+1)$ $\epsilon(i-1)$

 $A(\epsilon)$, namely all possible recurrence states of two randomly selected short sequences of N = 2, 3, and 4 consecutive points in a $K(K \gg N)$ length time series. Such microstates $A(\epsilon)$ are $N \times N$ small binary matrices and for a large enough randomly select number of sample M, the recurrence entropy S can be computed by.

$$S(\epsilon) = -\sum_{i=1}^{M} PA_i \ln PA_i$$

where PA_i is the probability of occurrence of a specific state $A_i(\epsilon)$.



$$u_{n+1}^{(i)} = (1-\epsilon)f(u_n^{(i)}) + \frac{\epsilon}{2} \left[f(u_n^{(i+1)}) + f(u_{tn}^{(i-1)}) \right]$$
(1)

where $u_n^{(i)}$ is a state variable at discrete time n = 0, 1, 2, ... and discrete position i = 1, 2, ...N in a onedimensional lattice of size N. The system at each spatial position undergoes a temporal dynamics described by a quadratic map

$$f(u) = u + \alpha u - \beta u^2 \tag{2}$$

where $0 \le u \le 1$. In the following we fix $\beta = 0.1$ and use α as a variable parameter. We will consider random initial conditions over the interval $0 \le u_0 \le 1$ and fixed boundary conditions: $u_n^1 = u_n^N = 0$.

Results



Figura 1: (a) Phase diagram of the coupled map lattice: it depicts the dominant temporal period of the spatio-tempora

Figura 3: (a-d) Time series of the maximal recurrence entropy max(S) computed for the defect signal (orange lines), superposed with the dynamical behavior of max(S) (blue lines) computed for stochastic time correlated signals following a power law spectrum $1/f^{\sigma}$, $\sigma = 2.1, 2.3, 2.5, 2.7$ respectively. Vertical lines are representative of the standard deviation of each windowed computed max(S). (e-h) shuffled version of all signal.



Figura 4: Dependence of the power-law exponent σ , of the correlated stochastic signal, with the nonlinearity parameter α .

attractor for given values of ε and α , for *beta* = 0.1 and *N* = 100; (b) The density of KS entropy as a function of the same parameters; (c) Space-time plot of the coupled quadratic map lattice (1) for $\varepsilon = 0.1$ and $\alpha = 2.85$.;



Figura 2: (a) Motion of an individual defect; (b) Number of chaotic defects as a function of the lattice size; (c) Mean square displacement of the chaotic defects as a function of time for a lattice of N = 6000 maps with $\alpha = 2.85$, $\beta = \varepsilon = 0.1$ (d) Diffusion coefficient of defects as a function of the nonlinearity parameter α

The recurrence plot is a tool to depict recurrences of a K length time series and is defined as a $K \times K$ binary matrix.

$$\mathbf{R}_{(i,j)}^{m,\epsilon_i} = \Theta(\epsilon - \|\mathbf{x}_i - \mathbf{x}_j\|), \mathbf{x}_i \in \mathbb{R}^m, i, j = 1, \dots, K$$
(3)

where ϵ is the vicinity parameter.

The concept of subsets of the recurrence plot has been generalized by using recurrence microstates

Conclusions

- we have revisited the formation and dynamics of chaotic defects in lattices of coupled quadratic maps. Chaotic defects are spatially localized regions of chaotic dynamics which propagate along the lattice in a particle-like fashion. Since we have used random initial conditions in numerical simulations we believe that the formation of chaotic defects is a quite general phenomenon.
- We found that the number of defects increases with the lattice size, and that defects experience a normal (gaussian) diffusion with a diffusion coefficient around 1.25. This value was found to depend on the nonlineary parameter Îs of the coupled maps.
- The motion of individual defects is similar to a random walk, and we used a recently developed recurrence-based technique to show that it is actually a stochastic time series displaying a power-law power spectrum $1/f^{\sigma}$, where σ lies between 2.3 and 2.4. The latter exponent was also found to be dependent on the nonlinearity parameter α . In other words, the defect motion has recurrence properties similar to a stochastic signal with memory.

Referências

- K. Kaneko and I. Tsuda, Complex Systems: Chaos and Beyond (Springer, New York-Berlin-Heidelberg, 2001)
- [2] N. Marwan, M. C. Romano, M. Thiel, and J. Kurths, Phys. Rep. 438, 237 (2007).
- [3] S. G. Corso, T. de L. Prado, G. Z. dos Santos, J. Kurths, S. R. Lopes, Chaos 28, 083108 (2018)
 [4] K. Kaneko, Physica D 34, 1 (1989)
- [5] R. Badii and A. Politi, Complexity, Hierarchical Structures and Scaling in Physics (Cambridge University Press, 1997)